Intermolecular bond stability of C_{60} dimers and 2D pressure-polymerized C_{60}

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The thermal stability of C_{60} dimers and 2D pressure-polymerized C_{60} is studied using high-resolution capacitance dilatometry. The transformation of both the dimer and the polymer phases back to 'normal' C_{60} is excellently described by a simple thermally activated process, with activation energies of 1.75 ± 0.1 eV (dimer) and 1.9 ± 0.2 eV (polymer). These results are compared to previous data of 1D-polymerized C_{60} and photo-polymerized C_{60} . The thermal expansivity of the 2D-polymer phase is as much as a factor of ten smaller than that of pure C_{60} and approaches values for diamond.

INTRODUCTION

Besides the 'normal', face-centered-cubic (fcc), plastic-crystal phase of C_{60} , several other phases have been reported in which the C_{60} molecules are linked via covalent bonds to form dimers, linear chains (1D), 2- and even 3-dimensional networks [1]. These polymerized phases can be produced by different methods, such as photopolymerization (dimers and 1D) [2], high-pressure-temperature conditions (1D, 2D and 3D) [3,4] and the recently published solid-state mechano-chemical reaction of C_{60} with KCN (dimers) [5]. The resulting polymer phases are metastable under ambient conditions and revert back to fcc C_{60} during heating to higher temperatures.

Here we study the kinetics of the polymer to fcc transformation of C_{60} dimers and 2D pressure-polymerized C_{60} using high-resolution dilatometry up to 500 K.

EXPERIMENTAL

The C_{60} dimers were synthesized by a solid-state mechano-chemical reaction of C_{60} with potassium cyanide [5]. The 2D-polymerized sample was prepared by annealing high-purity sublimed polycrystalline C_{60} at 830 K for 5 hours at a pressure of 2.0 GPa and then cooling the sample before releasing the pressure. Two high-resolution

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capacitance dilatometers with temperature ranges of 4-300 K [6] and 150-500 K [7], respectively, were used to measure the thermal expansion. Data were taken at constant heating (cooling) rates, and He exchange gas (10 mbar) was used to thermally couple the samples to the dilatometers. To characterize the dimer phase, Raman-spectra were taken with a FT-Raman spectrometer using a Nd:YAG-laser with a wavelength of 1064 nm.

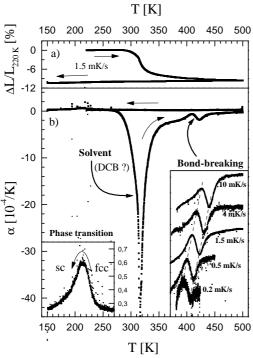


Fig. 1. Linear thermal expansion a) and expansivity b) of the C_{60} dimer.

RESULTS AND DISCUSSION

 C_{60} -dimer: In Fig. 1 we present the linear thermal expansion, $\Delta L/L_0$, and the expansivity, $\alpha = 1/L \cdot dL/dT$, C₆₀ dimer powder for the first heating cycle from 220 K to 500 K and the subsequent cooling cycle from 500 K to 150 K. Anomalies in $\alpha(T)$ are observed at about 315 K and 415 K in the heating curve and are absent in the cooling curve, indicating irreversible processes. The fccsc phase transition near 250 K is recovered in the cooling curve, demonstrating that the dimers have been broken apart by heating to 500 K. In order to determine which of the anomalies upon heating is due to bond breaking, Raman spectra [8] of fresh samples were taken before heating, after annealing at 385 K for three hours and after slowly heating to 500 K.

These spectra clearly demonstrate that the 315 K peak is due to evaporating solvent (o-dichlorobenzene) and that bond breaking occurs around 415 K, in good agreement with differential scanning calorimetry results from Ref. 5.

The right inset of Fig. 1b shows that the bond-breaking anomaly shifts to lower temperatures with decreasing heating rates. Defining the bond-breaking temperatures $T_{bb} \equiv T_{bond\text{-breaking}}$ as either the maxima or minima of the anomalies allows one to plot the logarithm of the heating rate r versus $1/T_{bb}$, from which an activation energy is obtained. The activation energies are $E_a = 1.70 \pm 0.05$ eV (maxima) and $E_a = 1.79$ eV ± 0.05 eV (minima). To convert the heating rate to a bond breaking rate τ_{bb}^{-1} , the bond breaking was simulated with a simple model in which the fraction of dimers y obeys

the differential equation
$$\frac{dy}{dT} = \frac{y}{\tau_U(T) \cdot r}$$
 (1)

with an activated rate $\tau_{bb}^{-1}(T) = v_0 \cdot \exp(-E_a/k_BT)$. The only free parameter in this model is the attempt frequency, which could be determined to $v_0 = (2.6 \pm 3) \cdot 10^{17}$ Hz.

The resulting temperature dependent bond breaking rate of the dimer samples is in very good agreement with that of the photo-polymerized C_{60} films from Y. Wang et al. [9], which strongly suggests that these films consist primarily of dimers. (see Fig. 3)

2D pressure-polymerized C_{60} : The thermal expansivity of the 2D-polymerized sample measured along two orthogonal directions is shown in Fig. 2. A surprisingly large anisotropy is observed, which possibly could arise from an uniaxial pressure component during synthesis. The expansivity in both directions is significantly smaller than that of both 'normal' C_{60} [10] and 1D-polymerized C_{60} [11], as expected due to stronger covalent bonding between molecules. The values along the direction with the smaller expansivity approach the very small thermal expansivity of diamond [12].

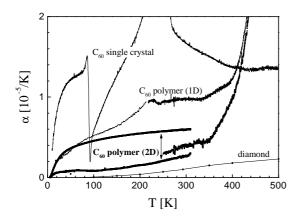


Fig. 2. Thermal expansivity of 2D-polymerized C_{60} compared to 'normal' C_{60} [10], 1D-polymerized C_{60} [11] and diamond [12].

The kinetics of depolymerization and the associated volume increase of the 2D-polymerized C₆₀ sample studied in detail by measuring the thermal expansion for repeated heating and cooling cycles between 150 K and 500 K (see Fig. 2 for the initial part of the first heating cycle). We obtained qualitatively very similar results as in our previous measurements on 1Dpolymerized C_{60} [11], but with two important differences. First, the length increase upon depolymerization of the 2D polymer was almost a factor of ten larger than for the 1D-polymer (6.3 % vs. 0.7 %), implying a quite large

volume increase in good agreement with the literature [3]. Second, although the bond-breaking rate τ_{bb}^{-1} exhibits an activated behavior with an activation energy $E_a = 1.9 \pm 0.2$ eV for the different heating and cooling cycles just as in the 1D-material, the attempt frequency decreases during the course of the measurement from an initial value close to the value of the 1D-polymer ($\nu_0 = 7 \cdot 10^{15}$ Hz) and ends up with $\nu_0 = 7.3 \cdot 10^{14}$ Hz (see Fig. 3). This may indicate that the sample is a mixture of 1D-and 2D-polymer phases, which are decomposing at different rates. A direct measurement of the time-dependent length-increase at 500 K near the end of depolymerization supports this interpretation; two different relaxation rates (see opencircles in Fig. 3), the values of which excellently match those of the 1D- and 2D-polymers obtained during heating and cooling cycles, are found.

Comparison: The bond-breaking rates of dimers, 1D- and 2D-polymerized materials are compared in Fig. 3. Interestingly, τ_{bb}^{-1} of the dimers is more than 3 orders of magnitude faster than that of the 1D or 2D material. This is quite surprising since experimental results [13] indicate that the dimers are energetically more stable than the

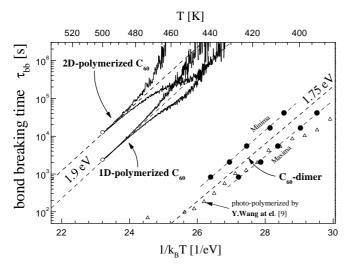


Fig. 3. Bond-breaking time τ_{bb} of the C_{60} dimer (\bullet) and the 2D-polymerized C_{60} compared to 1D-polymerized [11] and photopolymerized C_{60} (Δ) [9]. (see text for details)

1D or 2D materials. This difference appears to result both from different attempt frequencies and slightly different values of the activation energy. We note values the that the activation energies found here well agree very with theoretically determined values (~1.7 eV) for the 2Dpolymer phase [14].

CONCLUSIONS

The kinetics of bond-breaking of C_{60} dimers and 2D-polymerized C_{60} has been studied using high-resolution capacitance dilatometry up to 500 K. The bond-breaking rates are thermally activated with activation energies of 1.7 - 1.9 eV. The expansivity of the 2D-polymerized sample is much smaller than the expansivity of both 'normal' C_{60} [10] and even 1D-polymerized C_{60} [11], and approaches values of diamond [12].

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